Nonequilibrium Green's Function Simulation of Quantum Cascade Laser Structures

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Outline

- Intersubband Optics and Quantum Cascade Lasers
- Green's function modelling of QCLs
- Numerical Illustration
- Conclusion





Interband vs. Intersubband Optics:





interband transition

- •bipolar
- •photon energy determined by bandgap energy E_{gap} of material

intersubband transition

- •unipolar
- photon energy determined by well thickness, adjustable





The Challenge: Intersubband Lasing



Kazarinov and Suris, 1971





The Quantum Cascade Laser





LO-Phonon Assisted THz Design

S. Kumar et al.,

Appl. Phys. Lett., Vol. 84, No. 14, 5 April 2004







GREEN'S FUNCTION MODELLING OF QUANTUM CASCADE LASERS





Electron transport through quantum cascade lasers is determined by the balance between coherent transport and scattering:







Hamilton operator: $H = H^{(0)} + H_{\text{scatt}}$

Coherent transport governed by single-particle Hamiltonian: $H^{(0)} = H_{kin} + V_{SL} + V_{field}$

Scattering through many-body interactions in H_{scatt}:

- LO-phonon scattering
- acoustic phonon scattering (currently modelled by using an artifical LO-phonon dispersion)
- Impurity scattering
- Interface roughness scattering
- electron-electron scattering
- radiative losses





Both coherent transport and scattering can conveniently be described on the same footing using **Green's functions**:

$$G_{\alpha\beta,\mathbf{k}}^{<}(\omega) = \int dt \left\langle \Psi_{\beta,\mathbf{k}}^{\dagger}(t) \Psi_{\alpha,\mathbf{k}}(0) \right\rangle e^{i\omega t}$$

Relation to (single particle) density matrix:

$$\rho_{\alpha\beta,\mathbf{k}} = \int d\omega \ G^{<}_{\alpha\beta,\mathbf{k}}(\omega)$$

S.-C. LEE and A. WACKER, Phys. Rev. B 66, 245314 (2002). S.-C. LEE, F. BANIT, M. WOERNER, and A. WACKER, Phys. Rev. B 73, 245320 (2006).





α , β : Wannier-Stark states

(eigenstates of superlattice + external field)









Other Green functions:

$$G^{>}_{\alpha\beta,\mathbf{k}}(\omega) = -i \int dt \left\langle \Psi_{\alpha,\mathbf{k}}(0) \Psi^{\dagger}_{\beta,\mathbf{k}}(t) \right\rangle e^{i\omega t}$$

Spectral function:

$$\hat{G}_{\alpha\beta,\mathbf{k}}(\omega) = i[G^{>}_{\alpha\beta,\mathbf{k}}(\omega) - G^{<}_{\alpha\beta,\mathbf{k}}(\omega)]$$

Retarded Green function (Lehmann representation):

$$G_{\alpha\beta,\mathbf{k}}^{\text{ret}}(\omega) = \frac{1}{2\pi} \int d\omega' \, \frac{\hat{G}_{\alpha\beta,\mathbf{k}}(\omega')}{\omega - \omega' + i\varepsilon}$$

 $(\Rightarrow \hat{G}_{\alpha\beta,\mathbf{k}}(\omega) = -2 \operatorname{Im} G_{\alpha\beta,\mathbf{k}}^{\operatorname{ret}}(\omega)$ if \hat{G} is purely imaginary)





(retarded) Dyson equation

$$G_{\alpha\beta,\mathbf{k}}^{\mathrm{ret}}(\omega) = \left((G_0^{\mathrm{ret}})_{\mathbf{k}}^{-1}(\omega) - \Sigma_{\mathbf{k}}^{\mathrm{ret}}(\omega) \right)_{\alpha\beta}^{-1}$$

 $(G_0^{\text{ret}})^{-1}_{\alpha\beta,\mathbf{k}}(\omega) = \hbar\omega \cdot \delta_{\alpha\beta} - H_{\alpha\beta}^{(0)} + i\varepsilon, \quad H^{(0)} = H_{\text{kin}} + V_{\text{SI}} + V_{\text{field}}$

Keldysh relation (Dyson equation for G^{\gtrless})

$$G_{\alpha\beta,\mathbf{k}}^{\gtrless}(\omega) = \sum_{\alpha'\beta'} G_{\alpha\alpha',\mathbf{k}}^{\text{ret}}(\omega) \cdot \Sigma_{\alpha'\beta',\mathbf{k}}^{\gtrless}(\omega) \cdot \left(G_{\beta\beta',\mathbf{k}}^{\text{ret}}(\omega)\right)^{\bullet} + initial \ condition \ term$$

 \rightarrow G as functional of Σ !





Illustration: only one state, constant damping $\Sigma^{ret}(\omega) = i\Gamma$

$$G_{\mathbf{k}}^{\mathrm{ret}}(\boldsymbol{\omega}) = \frac{1}{\boldsymbol{\omega} - E_{\mathbf{k}} + i\Gamma}$$

Spectral function

$$\hat{G}_{\mathbf{k}}(\omega) = \frac{1}{(\hbar\omega - E_{\mathbf{k}})^2 + \Gamma^2}$$

Electron distribution in thermal equilibrium

$$G_{\mathbf{k}}^{<}(\omega) = i f(\hbar \omega) \frac{1}{(\hbar \omega - E_{\mathbf{k}})^{2} + \Gamma^{2}}$$





LO phonon self-energy (second Born approximation)

$$\Sigma_{\alpha\beta,\mathbf{k}}^{\gtrless}(\omega) = \sum_{\substack{\alpha'\beta'\\\mathbf{k}',q_{z},\pm}} \pm b^{\gtrless}(\pm\hbar\Omega_{\mathrm{LO}}) \frac{M_{\alpha\alpha'}(-q_{z})M_{\beta'\beta}(q_{z})}{(\mathbf{k}-\mathbf{k}')^{2}+q_{z}^{2}} \cdot G_{\alpha'\beta',\mathbf{k}'}^{\gtrless}(\omega\mp\Omega_{\mathrm{LO}})$$

$$M_{\alpha\beta}(q_z) = \sqrt{\frac{\hbar\Omega_{\rm LO}e^2}{4\pi} \left(\frac{1}{e_{\rm so}} - \frac{1}{e_{\rm s}}\right)} \int dz \ e^{-iq_z z} \psi_{\alpha}^*(z) \psi_{\beta}(z)$$

Lehmann representation for Σ :

$$\Sigma_{\alpha\beta,\mathbf{k}}^{\text{ret}}(\omega) = \frac{1}{2\pi i} \int d\omega' \, \frac{\Sigma_{\alpha\beta,\mathbf{k}}^{>}(\omega') - \Sigma_{\alpha\beta,\mathbf{k}}^{<}(\omega')}{\omega - \omega' - i\varepsilon}$$

$\rightarrow \Sigma$ as functional of G \rightarrow self-consistent system of equations!





Approximation in current state-of-the-art simulator due to computational demand:

Evaluate matrix elements at 'typical wavevectors' ktyp, ktyp:

$$\Sigma_{\alpha\beta}^{\gtrless}(\omega) = \sum_{\substack{\alpha'\beta'\\\mathbf{k}',q_z,\pm}} \pm b^{\gtrless}(\pm\hbar\Omega_{\mathrm{LO}}) \frac{M_{\alpha\alpha'}(-q_z)M_{\beta'\beta}(q_z)}{(\mathbf{k}_{\mathrm{typ}} - \mathbf{k}'_{\mathrm{typ}})^2 - q_z^2} \cdot G_{\alpha'\beta',\mathbf{k}'}^{\gtrless}(\omega \mp \Omega_{\mathrm{LO}})$$

Choose \mathbf{k}_{typ} , \mathbf{k}'_{typ} (for $\alpha = \beta$):

$$\frac{\hbar^2 \mathbf{k}_{\rm typ}^2}{2m} = \hbar \Omega_{\rm LO}, \qquad \frac{\hbar^2 \mathbf{k}_{\rm typ}^{\prime 2}}{2m} = 2 \hbar \Omega_{\rm LO}$$







Self-Energy $\Sigma(\omega, k, \alpha, \beta)$







Self-Energy $\Sigma(\omega,k,5,5)$







Self-Energy $\Sigma(\omega,k,5,5)$















 $G < (\omega, k, 4, 4)$







 $G < (\omega, k, 3, 3)$







 $G < (\omega, k, 2, 2)$















Energetically and Spatially Resolved Density







Energetically and Spatially Resolved Current























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