



Simulation of EM wave propagating in a nanocylinder-base localized surface plasma resonance sensor

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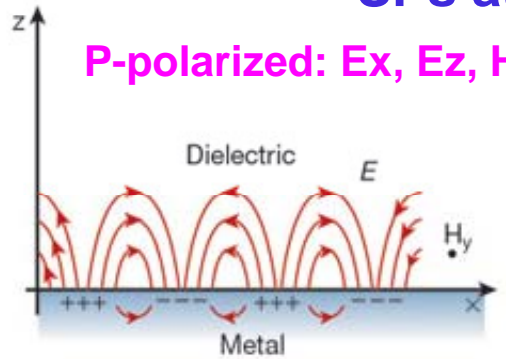
Outlines

- Introduction of surface plasmon resonance
- Experimental setup
- electronic resonance in single or pair of cylinders
- Field enhancement in Cylinder array
- Sensitivity enhancement
- Magnetic resonance
- conclusions

Introduction

Dispersion Relations of SPs

SPs at metal-dielectric interface

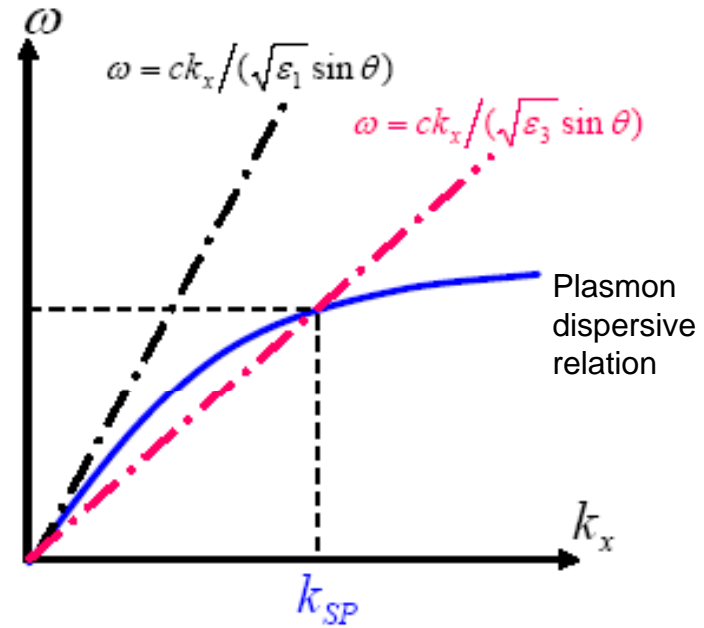
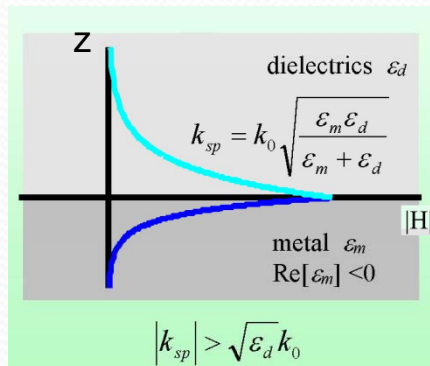


From Maxwell's equations & E and H fields continuity relations

→ Dispersion relation (DR)

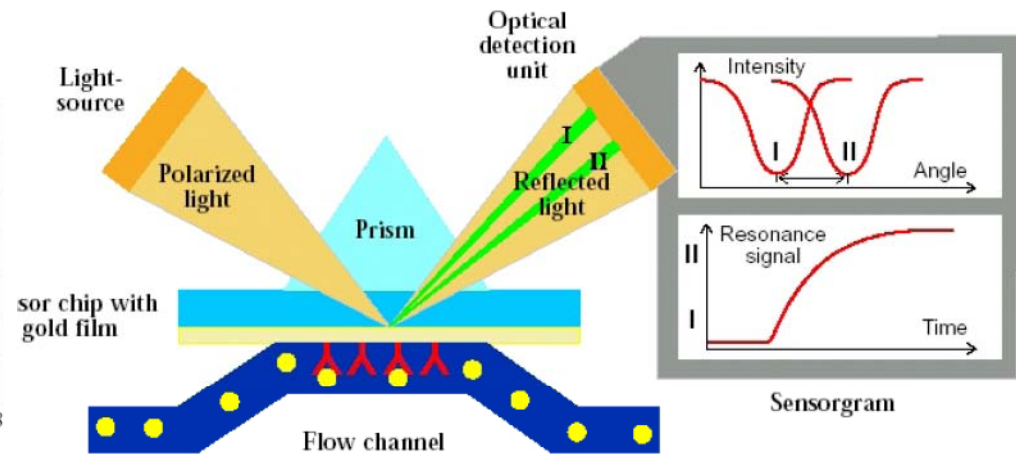
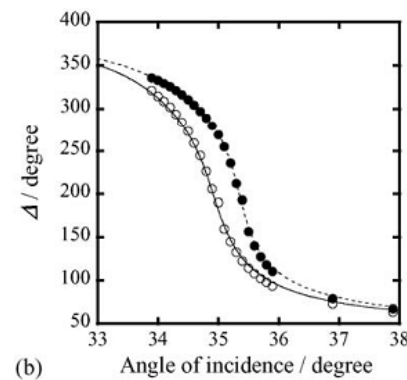
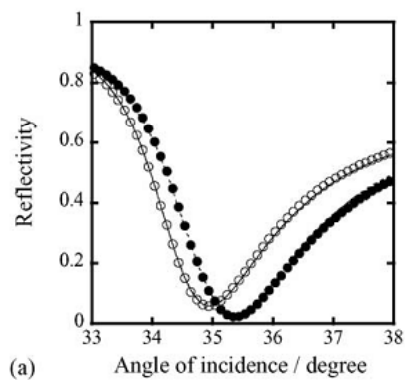
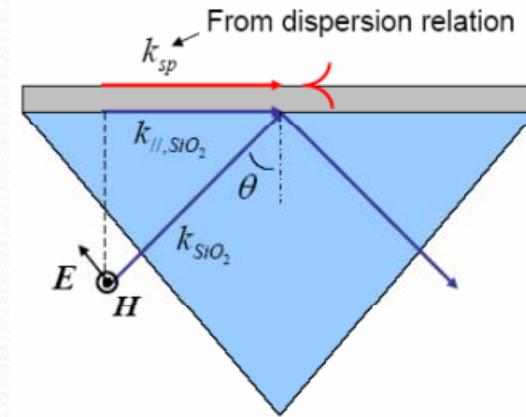
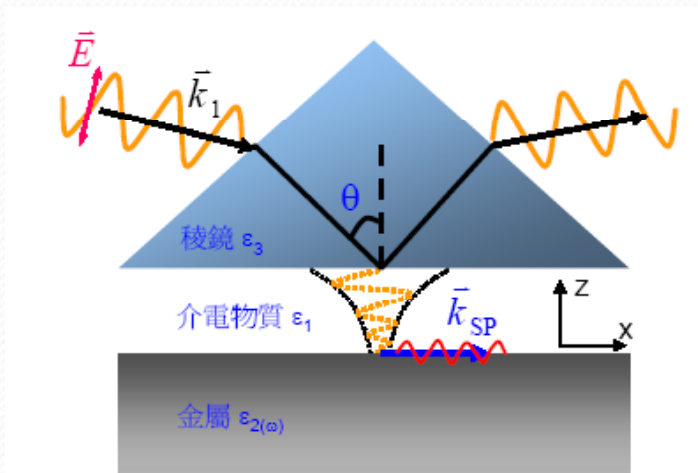
+ for $z \geq 0$, - for $z \leq 0$

k_z : imaginary



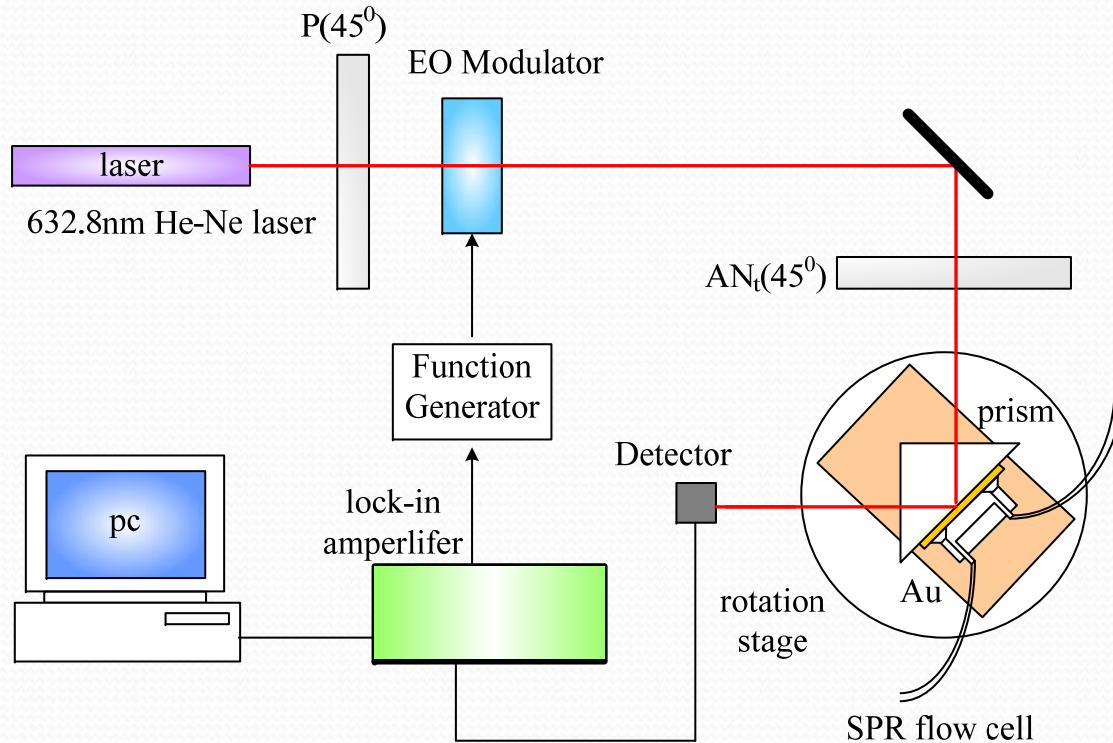
Excitation of SPs

ATR (attenuated total reflection) coupler



SP sensor

Experimental System



Jones vectors

$$P = \frac{1}{2} \begin{bmatrix} 1 \\ 1 \end{bmatrix} e^{i\omega_0 t}$$

$$\text{analyzer } AN_t = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$$

$$SPR = \begin{bmatrix} |r_p| e^{i\phi_p} & 0 \\ 0 & |r_s| e^{i\phi_s} \end{bmatrix}$$

$$EO = \begin{bmatrix} e^{i\frac{\omega t}{2}} & 0 \\ 0 & e^{-i\frac{\omega t}{2}} \end{bmatrix}$$

$$E_t = AN_t \cdot SPR \cdot EO \cdot P \quad I_t = E_t * E_t$$

$$I_t = \frac{1}{4} \left[\frac{|r_p|^2 + |r_s|^2}{2} + |r_p| |r_s| \cos(\omega t + \phi_p - \phi_s) \right]$$

$$\rho = \frac{r_p}{r_s} = \frac{E_{rp}/E_{ip}}{E_{rs}/E_{is}} = \tan \psi \exp(i\Delta)$$

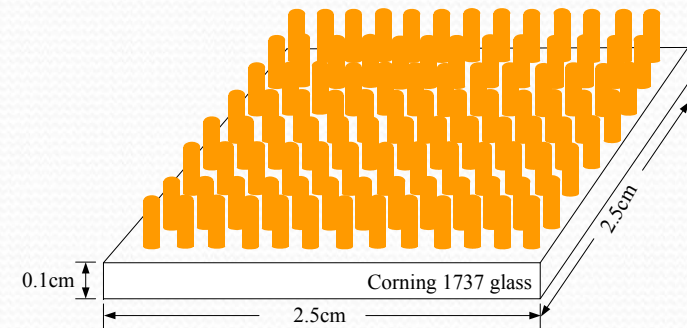
In general

$$\Rightarrow |r_s| = 1 \quad \phi_s \sim \text{constant}$$

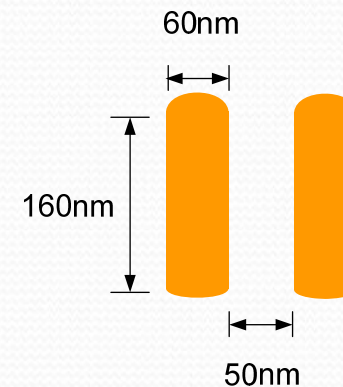
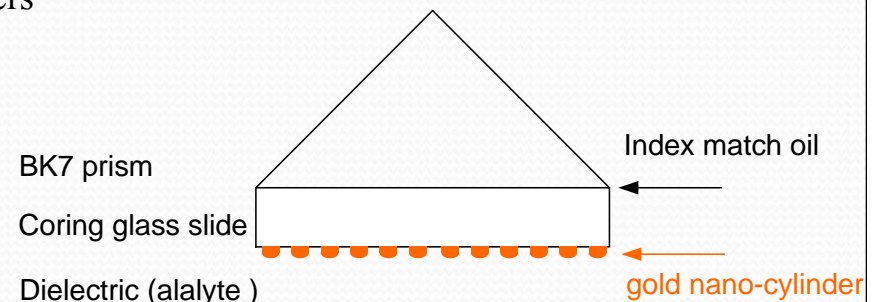
Mode

- Nanoscale cylinder array

- Interaction between localized and delocalized surface plasmaon polariton modes in a metal photonic crystal (A. Christ et al, *Phys. Stat. sol. (b)* 243. No.10 2344-2348)
- Plasmon resonance coupling in metallic nanocylinders (Kottmann et al, *Optics Express* Vol.8 No.12 P655 2001, Sburlan et al, *Physical Review B* 73 035403 2006)
- Nanowire-based enhancement of localized surface plasmon resonance for highly sensitive detection (K.Kim, *Optics Express* Vol.14, No.25, 12419 2006)
- Resonance light interaction with plasmonic nanocylinder system (ref: Viktor A. Podolskiy et al *J. of Optics A* 7 P32-37 2005)



Nano-cylinder array



Macroscopic permittivities

- Maxwell's equation in the metallic cylinder

$$\nabla \cdot \mathbf{D} = 4\pi\rho = \frac{4\pi}{i\omega} \nabla \cdot \mathbf{J} = \nabla \cdot \left(\frac{4\pi\sigma}{i\omega} \mathbf{E} \right) \Rightarrow \nabla \cdot (\varepsilon \mathbf{E}) = 0 \quad \varepsilon \equiv \varepsilon_\infty - \frac{4\pi\sigma}{i\omega}$$

$$\langle \mathbf{M} \rangle = \frac{1}{2cV} \int d^3r \theta_1(r) \left[\mathbf{r} \times \left(\mathbf{J}(r) - \frac{i\omega}{4\pi} \mathbf{D}(r) \right) \right] = -\frac{i\omega}{8\pi c} \varepsilon_1 \int d^3r \theta_1(r) [\mathbf{r} \times \mathbf{E}(r)]$$

$$\langle \mathbf{B} \rangle = \frac{c}{\omega} (\mathbf{k} \times \langle \mathbf{E} \rangle) = \sqrt{\varepsilon_e \mu_e} (\mathbf{e}_k \times \langle \mathbf{E} \rangle) \equiv \mu_e \langle \mathbf{H} \rangle = \mu_e \langle \mathbf{B} - 4\pi \mathbf{M} \rangle \quad \langle \varepsilon \mathbf{E} \rangle \equiv \varepsilon_e \langle \mathbf{E} \rangle$$

$$\Rightarrow \langle \mathbf{M} \rangle = \frac{\mu_e - 1}{4\pi} \mathbf{B} = \frac{\mu_e - 1}{4\pi} \sqrt{\varepsilon_e \mu_e} (\mathbf{e}_k \times \langle \mathbf{E} \rangle) \quad \varepsilon_e \equiv \varepsilon_\infty \left(1 - \frac{\omega_p^2}{\omega^2 - i\omega\gamma} \right)$$

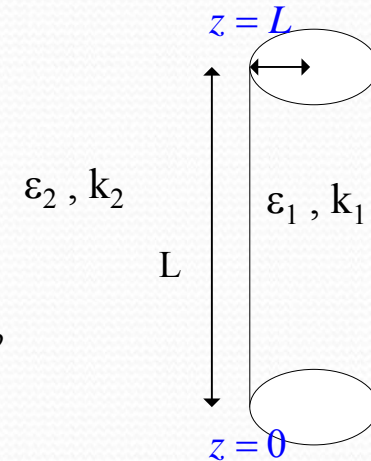
- Enhanced E-field make contribution ε_e , but μ_e limited.

Mathematic Description

- Single cylinder

- $L \gg a, \lambda > L$, the standing waves are good approximation to the eigenstates: $|M| \geq 1$

$$E_M(r) = -\cos(k_z z) \sum_{q=\pm 1} (q + \text{sign}M) \frac{k_z^2}{k_2^2} e_q J_{M-q}(\alpha\rho) e^{i(M-q)\varphi} + \sin k_z z \sqrt{2} \alpha \left(\frac{k_z}{k_2} \right) \text{sign}M e_0 J_M(\alpha\rho) e^{iM\varphi}$$



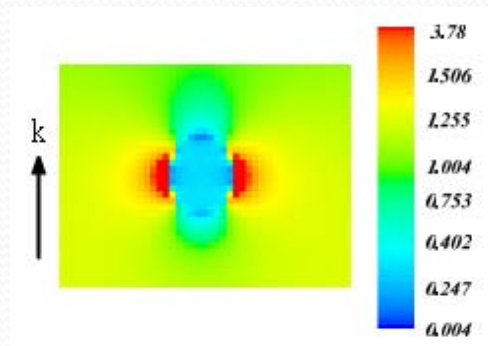
B.C.

$$E_M = 0 \text{ at } z = 0, L$$

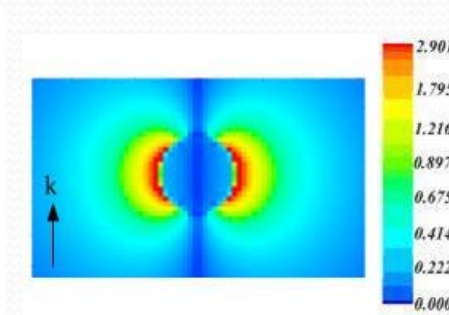
$$e_{\pm 1} \equiv \mp \frac{e_x \pm ie_y}{\sqrt{2}}, e_0 = e_z; k_i^2 \equiv \varepsilon_i \frac{\omega^2}{c^2} = \alpha_i^2 + k_z^2, i = 1, 2$$

$$k_z L = n\pi,$$

$$n = k_2 z / \pi, k_2 z / \pi + 1, \dots$$

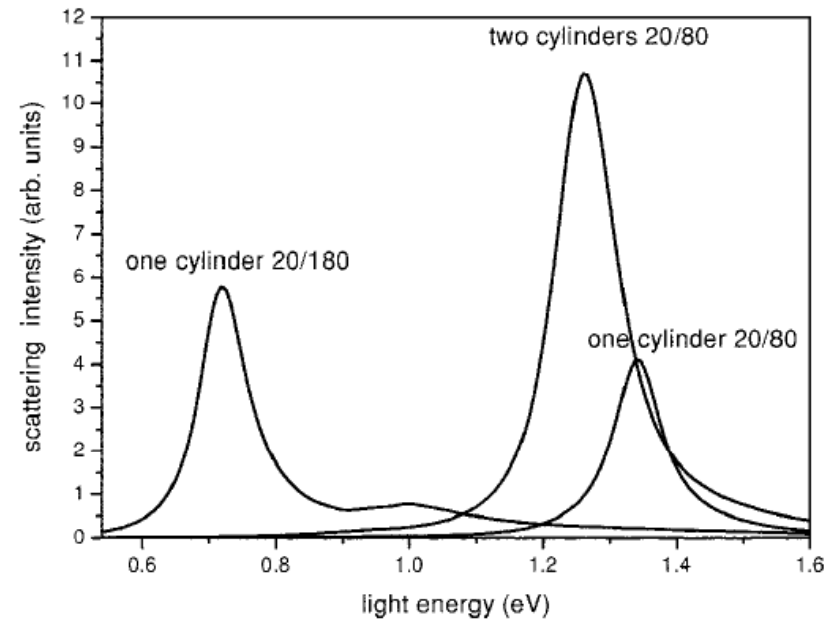
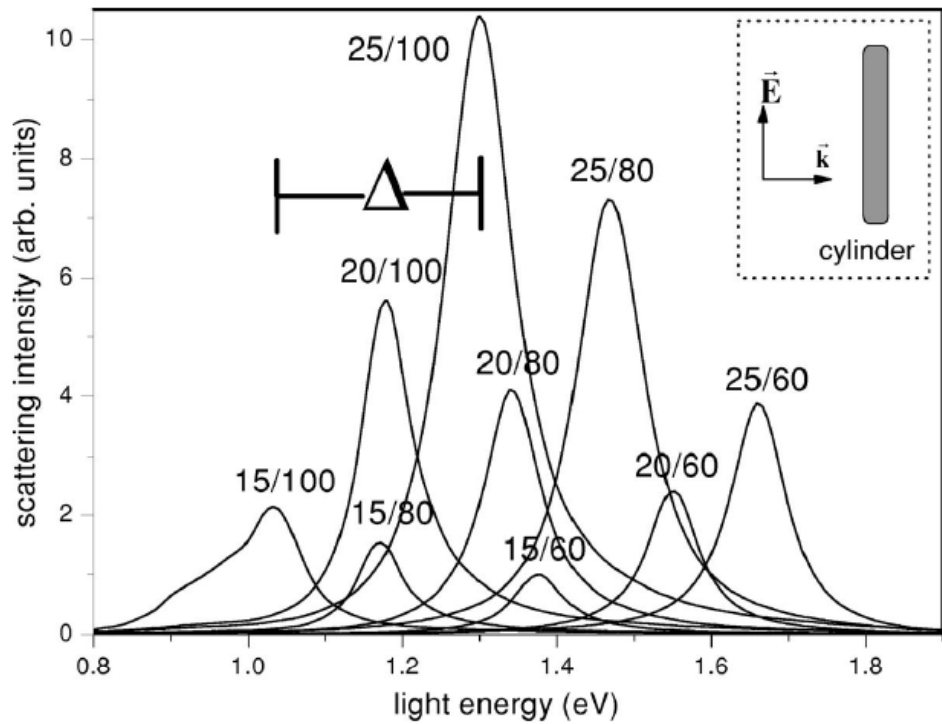


TM Mode



TE Mode

- TM mode (p-wave) field enhancement is a function of cylinder dimension



- Two cylinders

- $\rho < a$

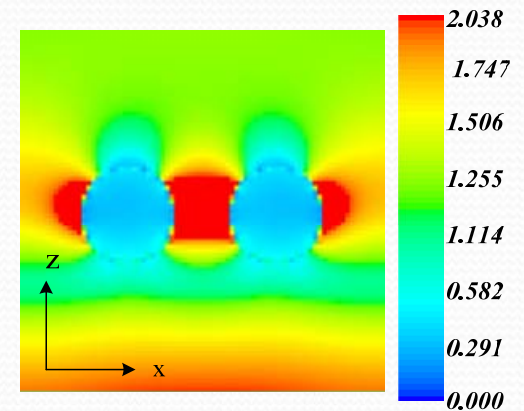
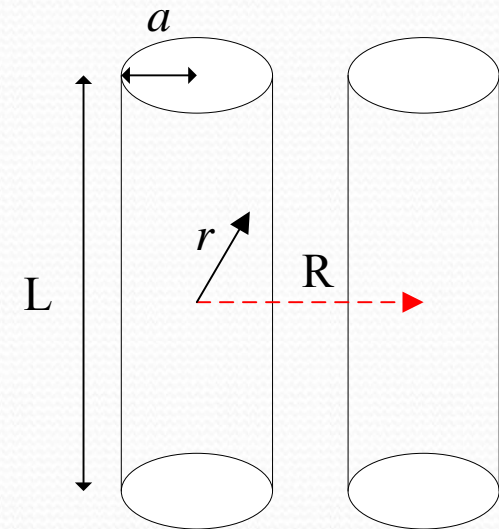
$$E_M(r) = -\cos(k_z z) \sum_{q=\pm 1} (q + \text{sign}M \frac{k_z^2}{k_2^2}) e_q J_{M-q}(\alpha_1 \rho) e^{i(M-q)\varphi}$$

$$+ \sin k_z z \sqrt{2} \alpha_1 (k_z / k_2^2) \text{sign}M e_z J_M(\alpha_1 \rho) e^{iM\varphi}$$

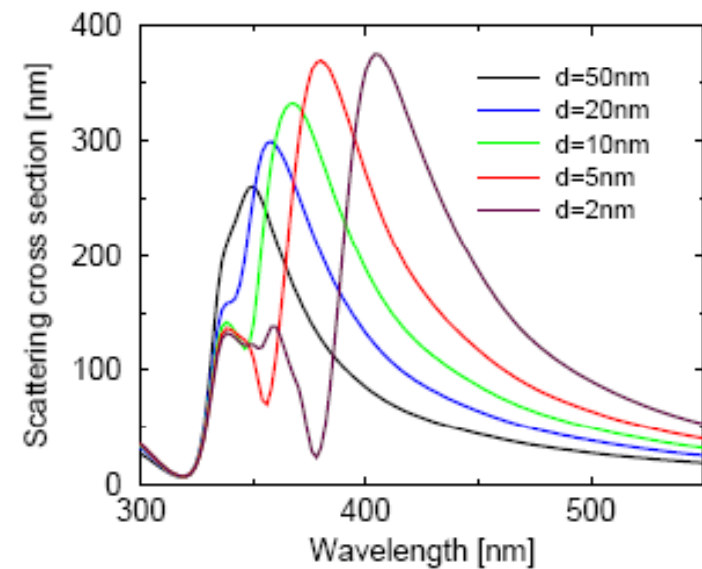
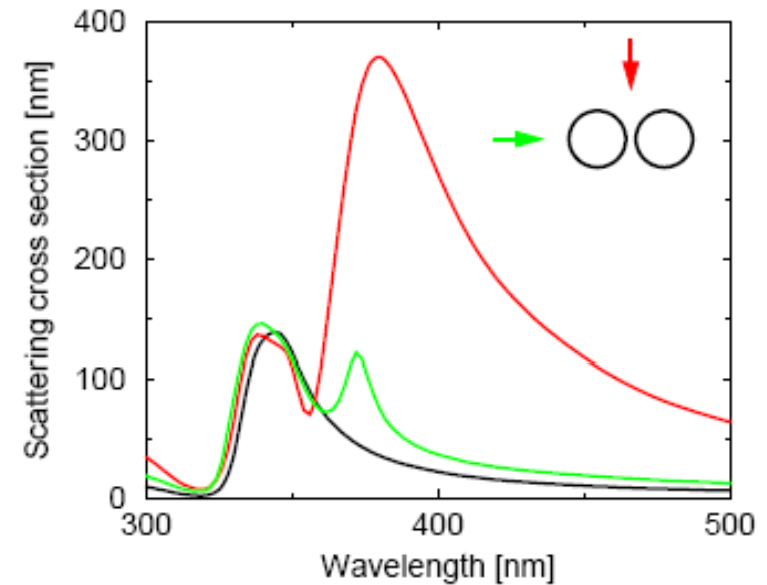
- $\rho > a$

$$E_M(r) = -\cos(k_z z) \sum_{q=\pm 1} (q + \text{sign}M \frac{k_z^2}{k_2^2}) e_q Y_{M-q}(\alpha_2 \rho) e^{i(M-q)\varphi}$$

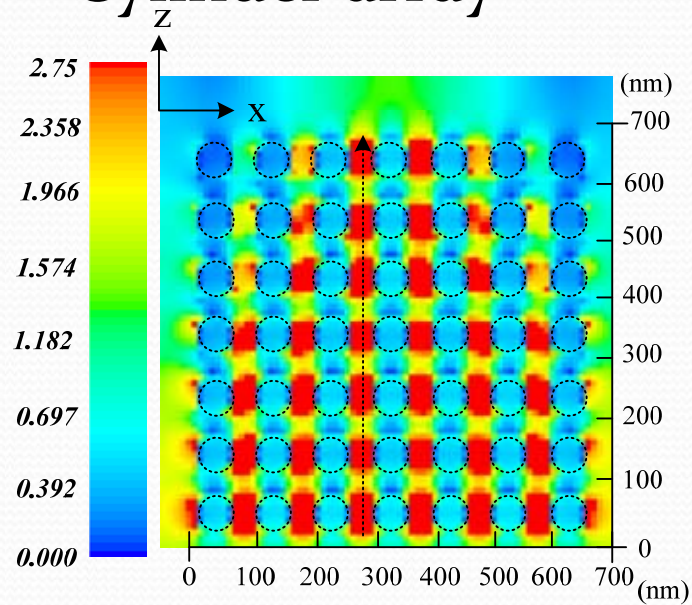
$$+ \sin k_z z \sqrt{2} \alpha_2 (k_z / k_2^2) \text{sign}M e_z Y_M(\alpha_2 \rho) e^{iM\varphi}$$



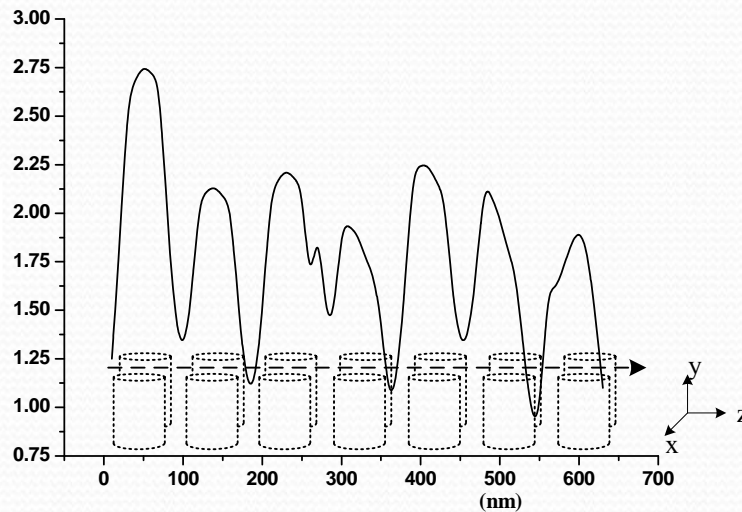
- Coupled system has the same magnitude resonance as for the individual cylinder, but the coupling is much stronger and depends on the incident field
- Main resonance is red-shifted with decreasing distance d



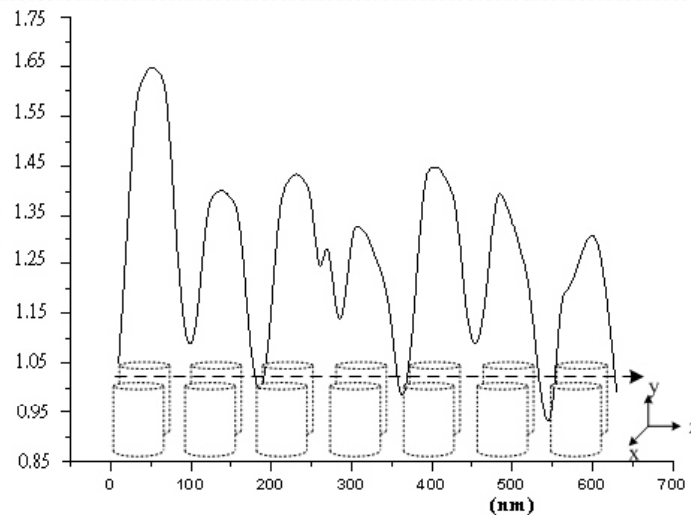
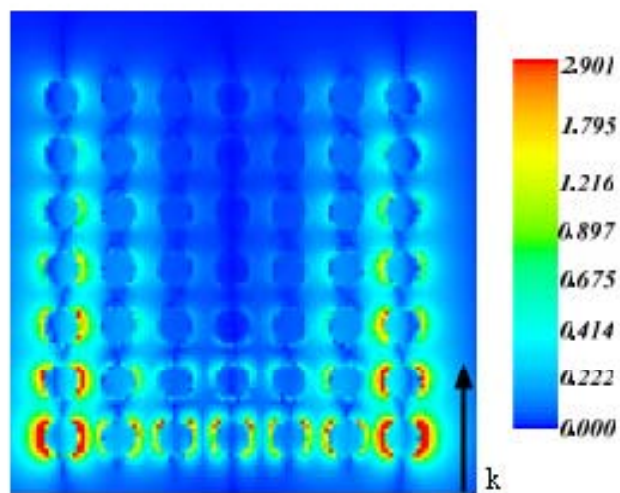
● Cylinder array



TM Mode



TE Mode



Local Surface Plasmon Enhance Sensitivity

- Sensitivity Definition: Mathematic formula

$$S = \frac{\Delta h}{\Delta n}$$

Δn : Refraction index variation

Δh : Physical parameter variation (θ or λ)

- Resolution

$$\sigma_{RI} = \frac{\sigma_{ins}}{S}$$

σ_{ins} : Instrument Resolution

- Present

σ_{RI} :

Can reach $\sim 10^{-7}$ for using heterodyne interferometer

Interrogation	Angular	Wave	Intensit y	Phase
σ_{ins}	0.0001°	0.02nm	0.2%	0.01°
Prism coupling	5×10^{-7}	2×10^{-5}	5×10^{-5}	4.6×10^{-7}
Grating coupling	2×10^{-6}	6×10^{-5}	2×10^{-4}	

J. Homola, S.S.Yee and G.Gauglitz "Surface plasma resonance sensors: Review" Sens. Actuators B 54, 3-15 (1999)

- Local surface plasma resonance is accompanied by the broadening in dispersion relation which increase the damping term γ *
- Surface plasmon resonance consists of gold thin film part and nanocylinder part, angular width is

$$\Delta\theta_{sp} = \frac{\text{Im}\{k_{sp}\}}{n_s \frac{\omega}{c} \cos\theta_{sp}} = \frac{\gamma_i + \gamma_r}{n_s \frac{\omega}{c} \cos\theta_{sp}}$$

$$\epsilon_e \equiv \epsilon_\infty \left(1 - \frac{\omega_p^2}{\omega^2 - i\omega\gamma} \right) = \epsilon'_{ca} + i\epsilon''_{ca}$$

$$\text{Im}\{k_{sp}\} \approx \frac{\omega}{c} \left(\frac{\epsilon'_{Au} \epsilon'_{ca}}{\epsilon'_{Au} + \epsilon'_{ca}} \right) \frac{\epsilon''_{Au}}{2(\epsilon'_{Au})^2}$$

$$|\epsilon'_{Au}| \gg \epsilon''_{Au}$$

$$\epsilon''_{ca} \approx 0$$

Metallic photonic crystal

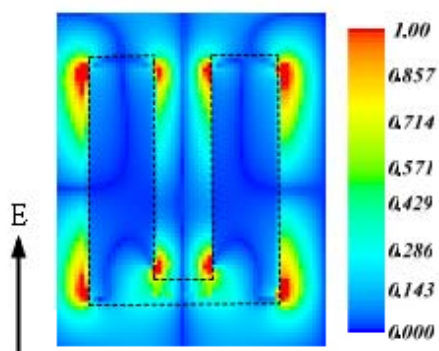
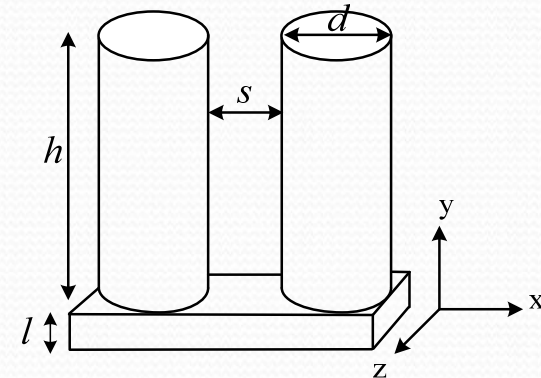
$$\Delta\theta_{sp} = \frac{2}{n_s \cos\theta_{sp}} \left(\frac{\epsilon'_{Au} \epsilon'_{ca}}{\epsilon'_{Au} + \epsilon'_{ca}} \right) \frac{\epsilon''_{Au}}{2(\epsilon'_{Au})^2}$$

$$|\epsilon'_{Au}| \sim \epsilon'_{ca}$$

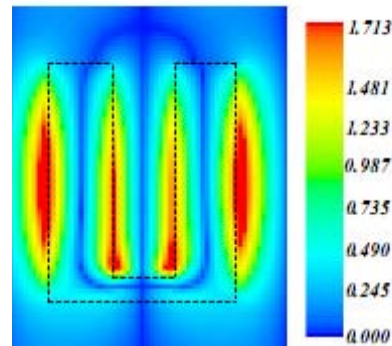
* Kyujung Kim Optics Express Vol 14, No 25 12419 (2006)

Magnetic Resonance

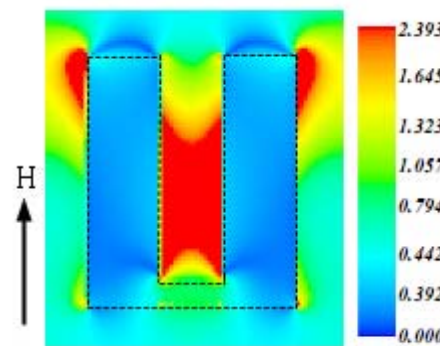
- TE Mode:
 - Anti-symmetric polariton mode
 - Anti-parallel current in the two cylinders induce magnetic dipole moment resonance
- TM Mode: Symmetric polariton will induce the electric resonance, but magnetic resonance is limited



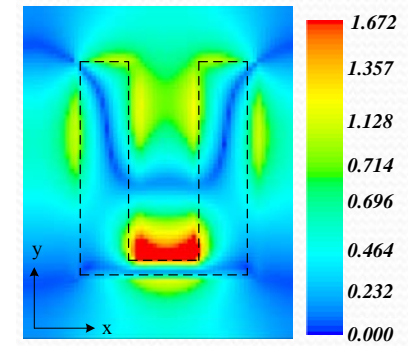
E in TE mode



H in TE mode



E in TM mode



H in TM mode

Conclusions

- Significant sensitivity increase is associated with larger field enhance.
- Field enhancement is achieved by nanocylinder based localized plasmon resonance mediated by nanocylinder and their coupling.
- Coupling has more contribution on the field enhance.
- Significant magnetic field enhancement is happen in TE mode due to anti-symmetric polarition mode.