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Dynamic Simulation of High Brightness Semiconductor Lasers

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Outline

- Modelling
- Numerical Solution
- Results
- Summary

Mathematical Modelling and Numerical Simulation

- Semiconductor lasers are characterized by a huge amount of physical and geometrical parameters
- Time-space instabilities like
 - Pulsations
 - Self-focusing
 - Filamentation
 - Thermal Lensing

→ Restrictions to beam quality and wavelength stability.

- Mathematical modeling and computer simulations are needed for
 - preparing technological processes to choose optimal parameters
 - understanding experimental data
 - predicting new laser designs.

2d Optical Amplitude Equations

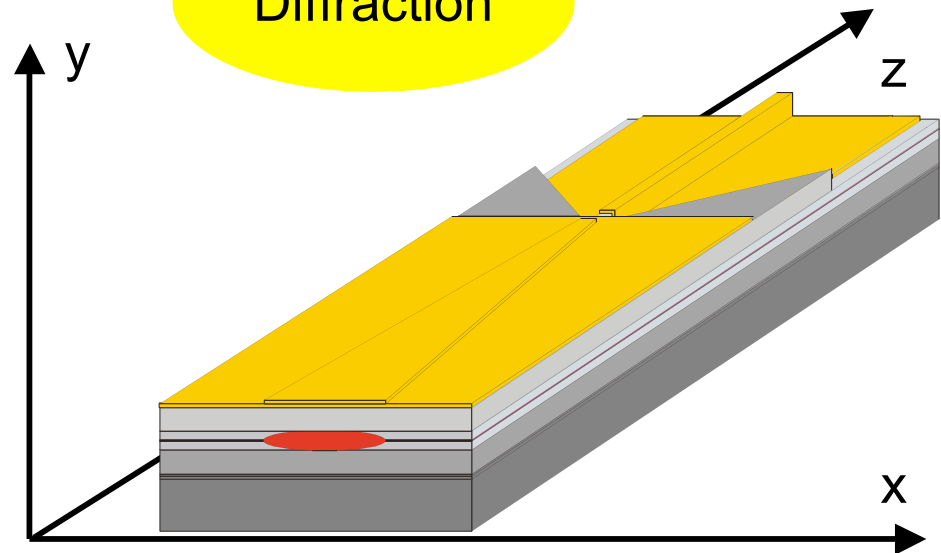
Propagation

$$\frac{1}{v_g} \partial_t u^\pm(x, z, t) = \mp \partial_z u^\pm(x, z, t) - i\beta(x, z) u^\pm(x, z, t) - \frac{1}{2\bar{n}k_0} i \partial_{xx} u^\pm(x, z, t)$$

Diffraction

with

$$\beta(x, z) := k_0 \frac{n_{eff}^2(x, z) - \bar{n}^2}{2\bar{n}}$$



Model for β

Effective refractive index depends on carrier density N and temperature.
We use the following model for β :

$$\beta = \delta_0(x, z) + \delta_{\Sigma}(x, z, N, J) + i \frac{g(x, z, N, u) - \alpha(x, z)}{2},$$
$$\delta_{\Sigma}(x, z, N, J) = \delta_n(x, z, N) + \delta_T(x, z, J),$$

where

- δ_0 index variation in space,
- $\delta_n = -\sqrt{n'N}$ dependence of effective index on the carrier density,
- δ_T temperature dependence of effective index,
- $g(x, z, N, u) = g'(x, z) \frac{\ln \frac{N(x, z, t)}{N_{tr}}}{1 + \epsilon \|u\|^2}$ optical gain,
- $\alpha = \alpha(x, z)$ optical losses.
- Model parameters determined at FBH and HU Berlin experimentally.

Model Equations

$$\begin{aligned}\frac{1}{v_g} \partial_t u^\pm &= \frac{-i}{2k_0 \bar{n}} \partial_{xx} u^\pm + (\mp \partial_z - i\beta) u^\pm - i\kappa u^\mp - \frac{\bar{g}}{2} (u^\pm - p^\pm) \\ \partial_t p^\pm &= \bar{\gamma} (u^\pm - p^\pm) + i\bar{\omega} p^\pm \\ \partial_t n &= d_n \partial_{xx} n + \frac{J}{qd} - R(n) - v_g \Re \langle u, g(n, u) u - \bar{g}(u - p) \rangle_{\mathbb{C}^2}\end{aligned}$$

$$u^+(t, 0, x) = r_0(x) u^-(t, 0, x) + \alpha(t, x), \quad u^-(t, l, x) = r_l(x) u^+(t, l, x).$$

- Spont. Recombination $R(n) = A(x, z)n + B(x, z)n^2 + C(x, z)n^3$
- Electrical injection $J = J(t, z, x)$, optical injection $\alpha(t, x)$
- Reflection coefficients $r_0(x), r_l(x)$
- All coefficients with the exception of k_0, v_g, \bar{n} are spatially, i.e. laterally and longitudinally in (z, x) plane, nonhomogeneous and discontinuous (depending on the heterostructural lasergeometry)

Heating

Considered parametrically via Injection $J(t, z, x)$:

$$I_{MO} = \int_{MO} J(x, z) dx dz, \quad I_{PA} = \int_{PA} J(x, z) dx dz.$$

Nonlocal dependence via:

$$\delta_T(x, z) = \frac{k_0 n_g}{\lambda_0} \int c_T(x, z, \tilde{x}, \tilde{z}) J(\tilde{x}, \tilde{z}) d\tilde{x} d\tilde{z},$$

Shift of gain peak:

$$\bar{\lambda}(x, z) = \bar{\lambda}_0(x, z) + \int \nu_T(x, z, \tilde{x}, \tilde{z}) J(\tilde{x}, \tilde{z}) d\tilde{x} d\tilde{z}.$$

Spectral method

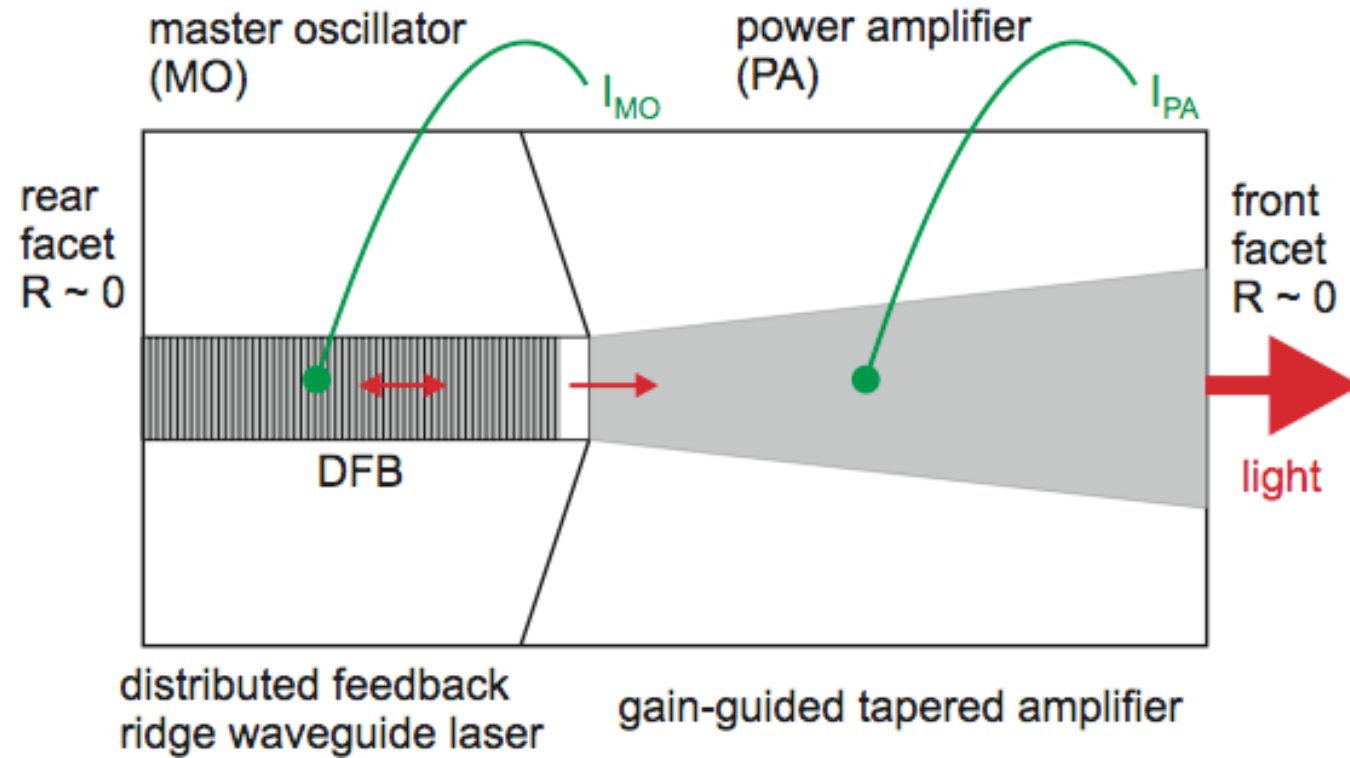
$$\frac{1}{v_g} \partial_t E^\pm = \frac{1}{2K} i \partial_{xx} E^\pm + (\mp \partial_z - i\beta(n)) E^\pm - i\kappa E^\mp - \frac{\bar{g}}{2} (E^\pm - P^\pm)$$

$$\partial_t P^\pm = \bar{\gamma} (E^\pm - P^\pm) + i\bar{\omega} P^\pm$$

$$\partial_t n = d_n \partial_{xx} n + I(t) - R(n) - v_g \Re \langle E, g(n) E - \bar{g} (E - P) \rangle_{\mathbb{C}^2}$$

- Step 1: Solve diffraction und diffusion (red coloring) with fast Fourier transform
- Step 2: Solve remaining 1d hyperbolic system for each lateral x coordinate via integration along characteristics for fields
 - Solve fast P equation using exact solution formula with forward value for fields
- For spectral split method good L^2 convergence has been proven for smooth scalar nonlinear Schrödinger equation
- Use predictor / corrector method for optical nonlinearities

MOPA



Large Scale Problem

- Longitudinal discretization $5\mu\text{m}$ for 973nm laser
- Time step 0.061ps
- Lateral discretization $0.625\mu\text{m}$
- Resulting # of spatial variables:
 - 2.88 million (4mm long MOPA)
 - 8 – 16mm laser: 5.76 – 11.52 million variables
- 1d parameter scan of different dynamical regimes requires $\approx 300\text{ns}$
- 2d scans 10–100 times more expensive

 Parallel computing required!

Remark:

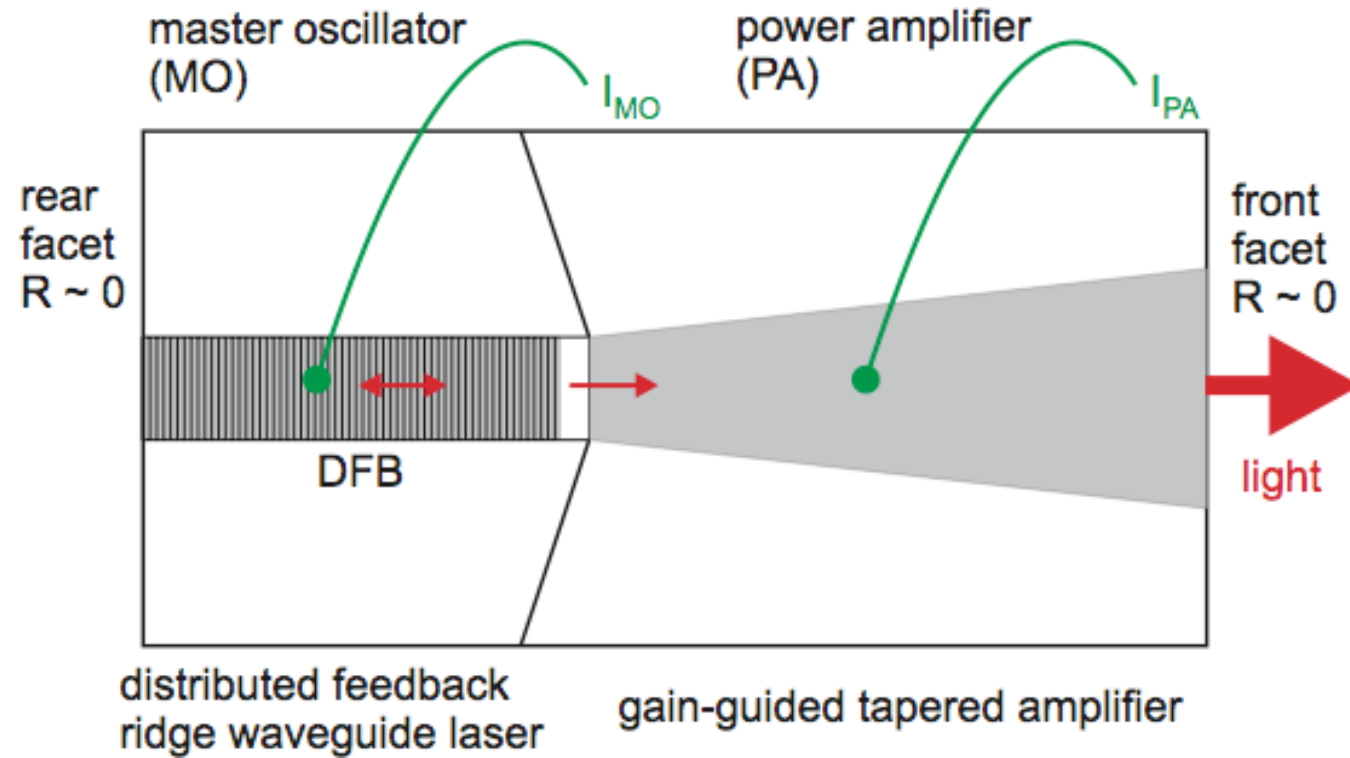
In comparison with previous 1d TWE model (*LDSL-tool*) calculation for HHI lasers our problem is 6400 times more complex!

High performance parallel computation on WIAS cluster Euler

10 × 300ns 2d parameter simulation on Euler cluster using multilayer parallel computing with MPI via Infiniband and POSIX Multithreading:

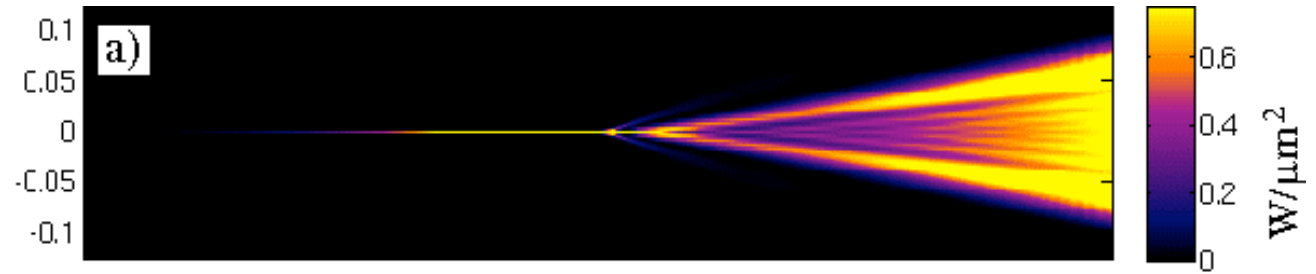
- Single PC: 2415 hours, 100 days
 - Euler 1 node (2 quad Xeons, 8 cores): 833 hours, 34 days
 - Euler 4 nodes (32 cores): 288 hours, 9 days
 - Euler 9 nodes (72 cores): 114 hours, 4.75 days
 - Euler 3×9 nodes (216 cores): 38 hours, < 2 days
 - total number of available nodes on Euler: 32 (256 cores)
- A single low resolution 2d parameter scan of dynamical regimes requires about 1 day on whole cluster.

Results

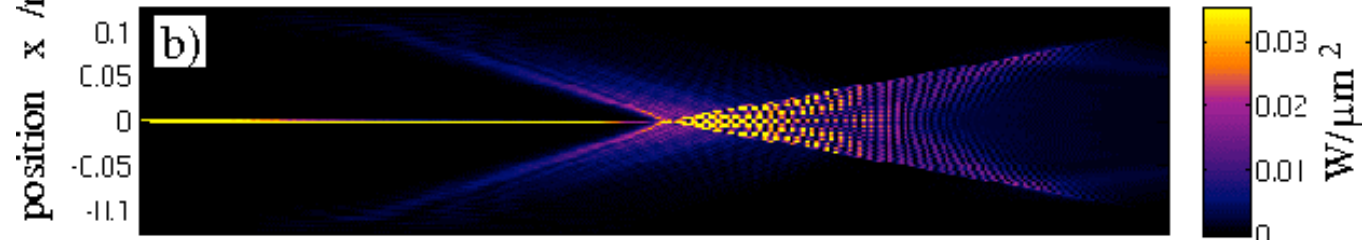


Intensity and Carrier Distributions

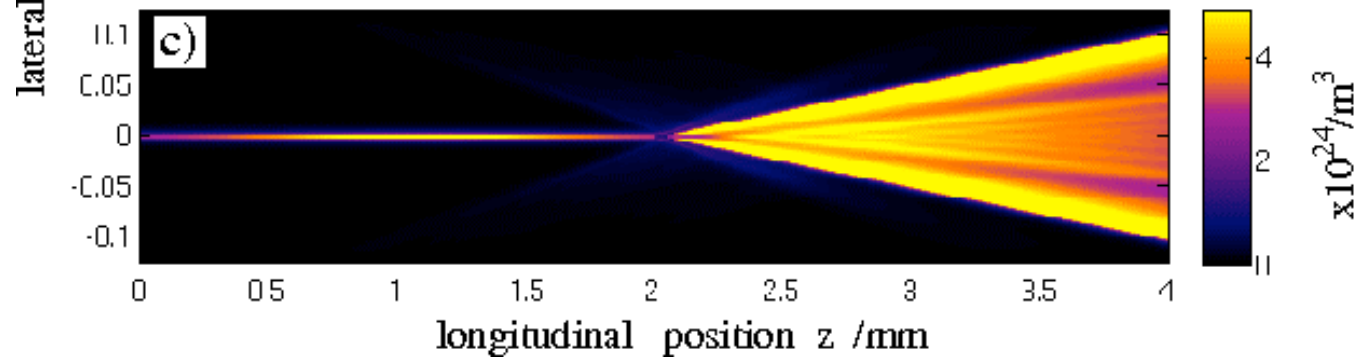
forward field:



backward field:



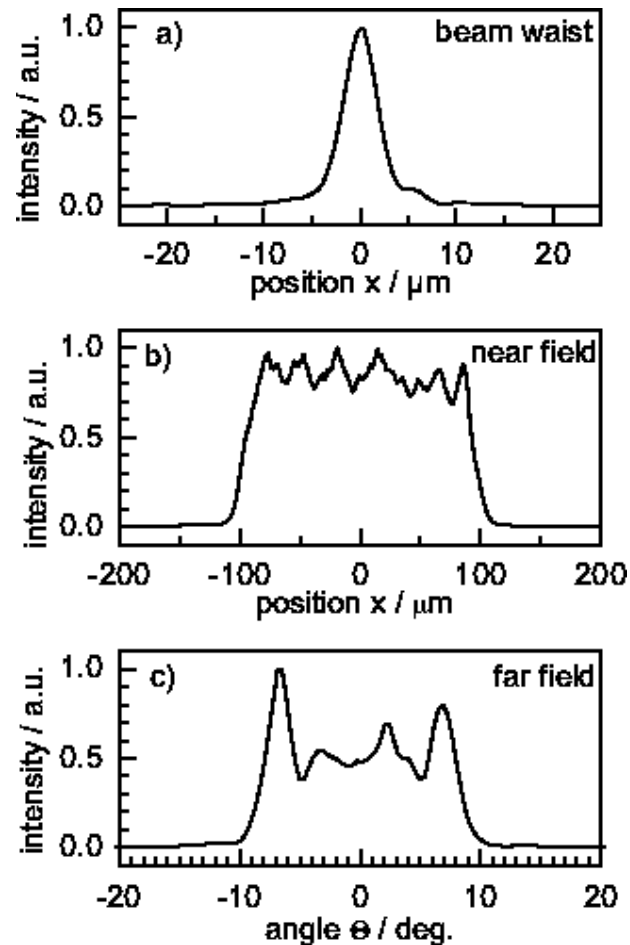
carrier density:



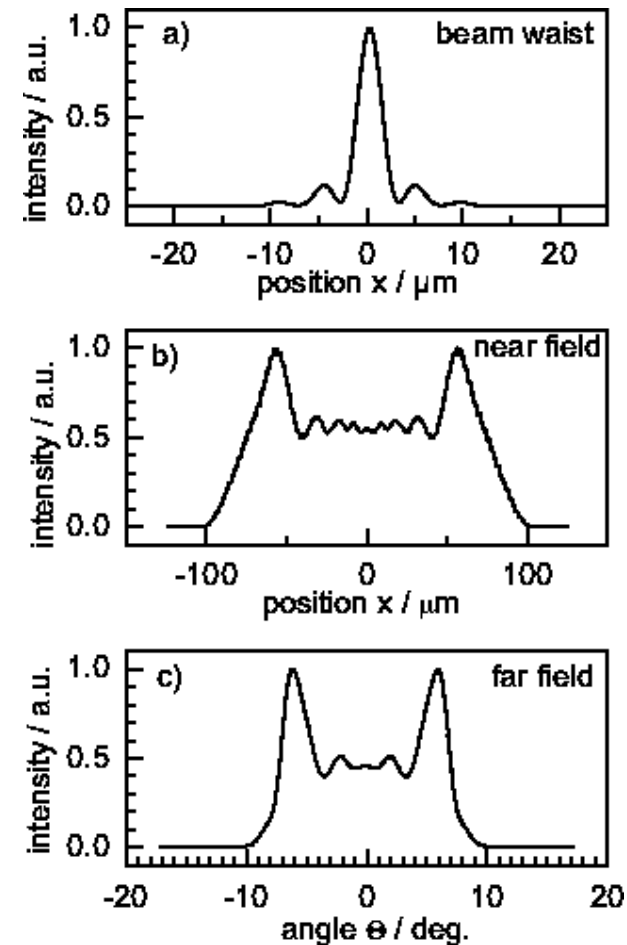
$$I_{MO}=400\text{mA}, I_{PA}=4 \text{ A}$$

Lateral Intensity Profiles

Experiment



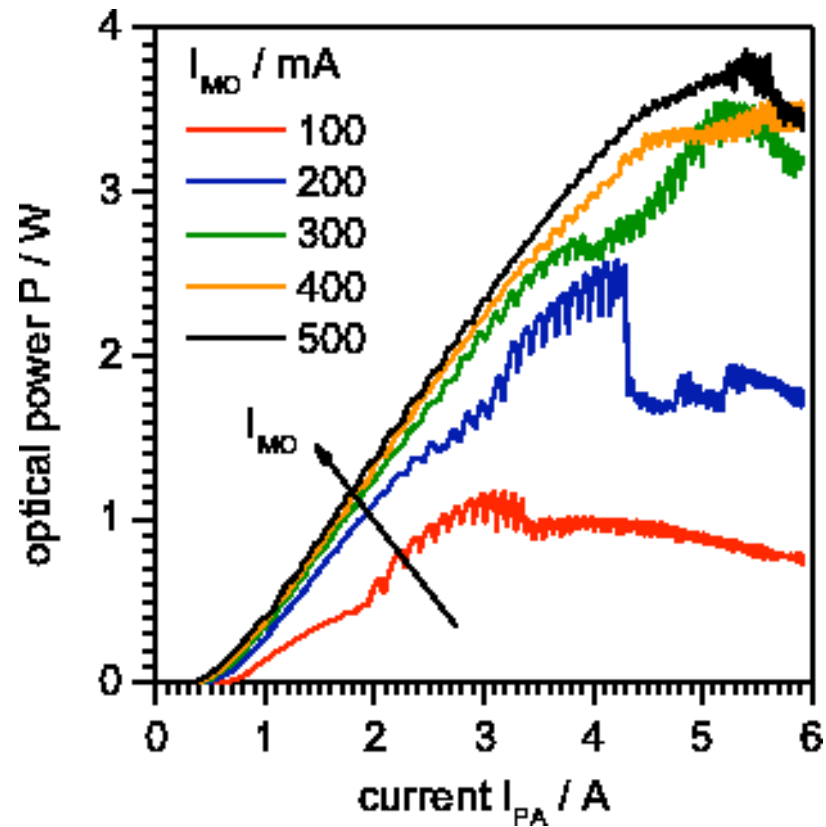
Simulation



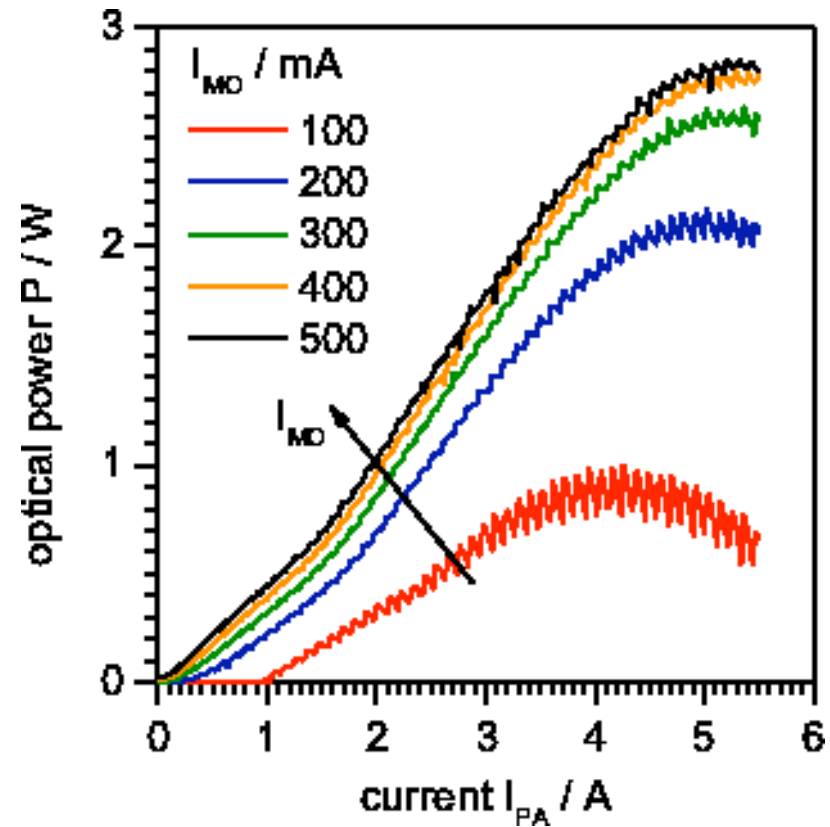
$$I_{MO}=400\text{mA}, I_{PA}=4 \text{ A}$$

P-I Characteristics

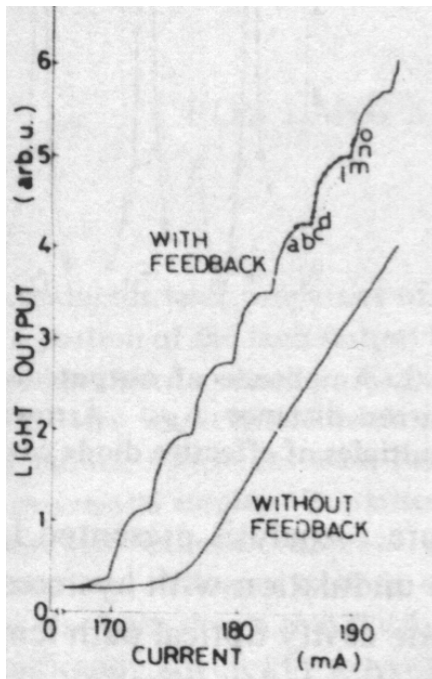
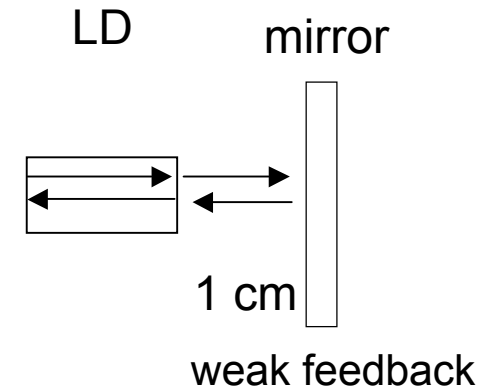
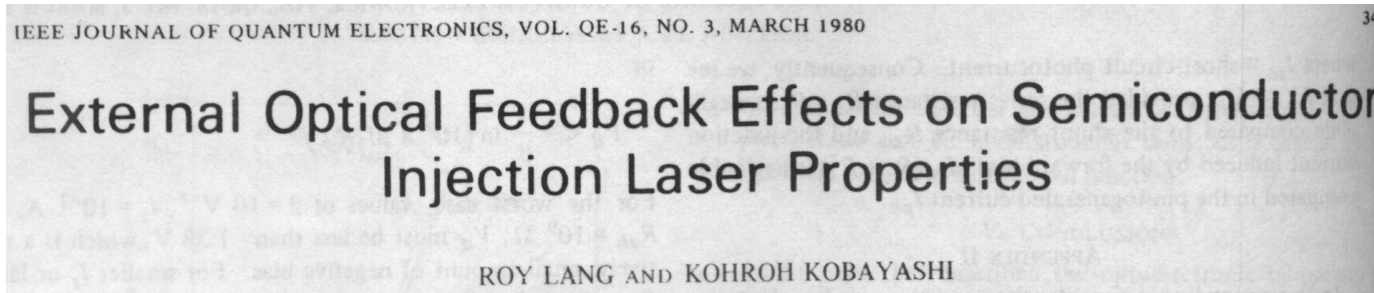
Experiment:



Theory:



Comparison with feedback experiments

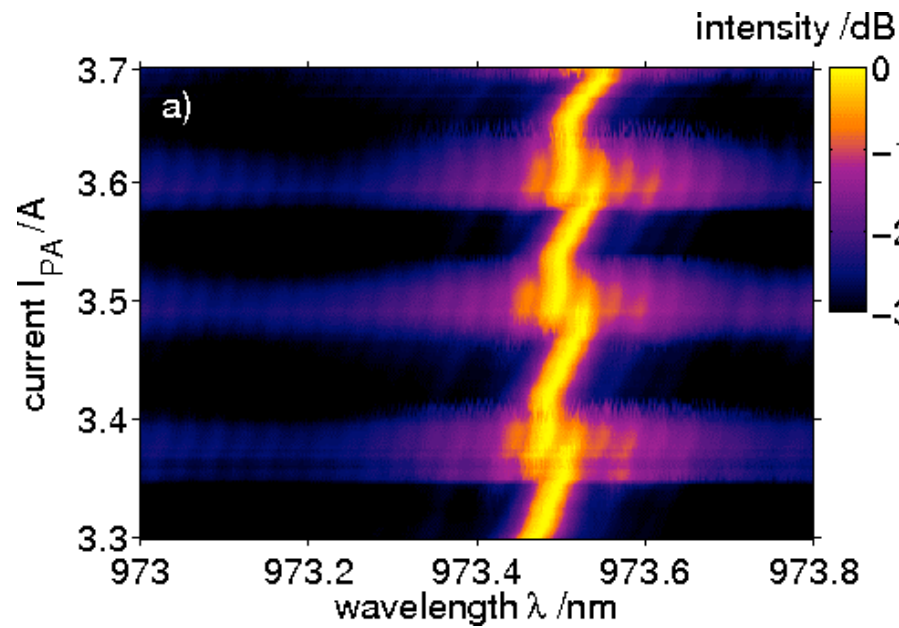


- laser + external mirror form compound cavity laser
 - coherent interference effects between reflected light and field inside the laser diode
 - $I \uparrow \Rightarrow \lambda \uparrow \Rightarrow$ feedback phase changes
 \Rightarrow intensity undulations
- Feedback from taper into DFB-laser !

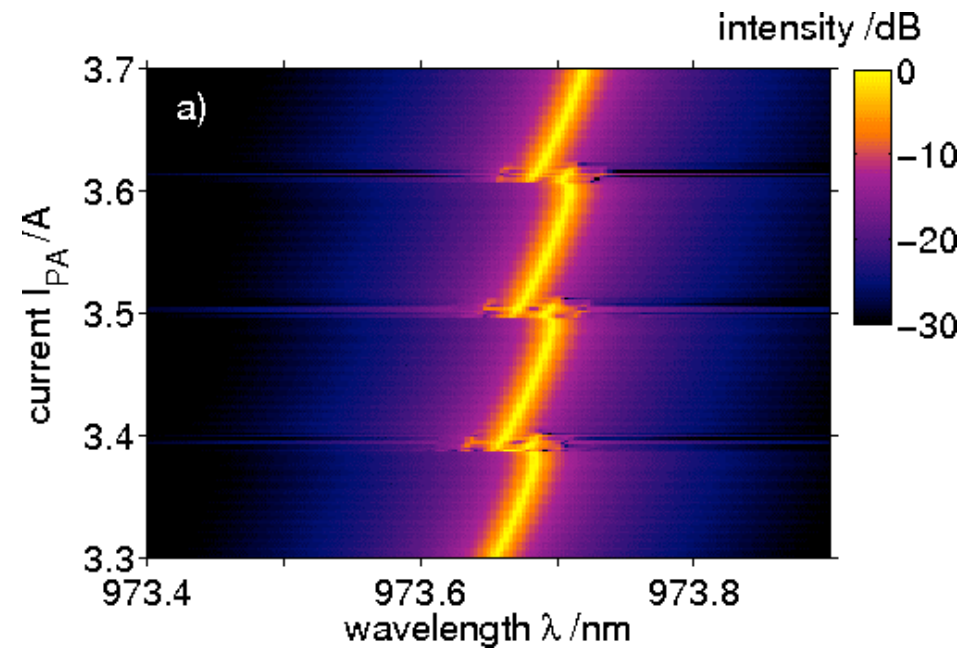
More recent for AFL: *S. Bauer et al., Phys. Rev. E 69 (2004) 016206*
O. Ushakov et al., Phys. Rev. Lett. 95 (2005) 123903

Optical Spectra vs I_{PA} :

Experiment:



Simulation:

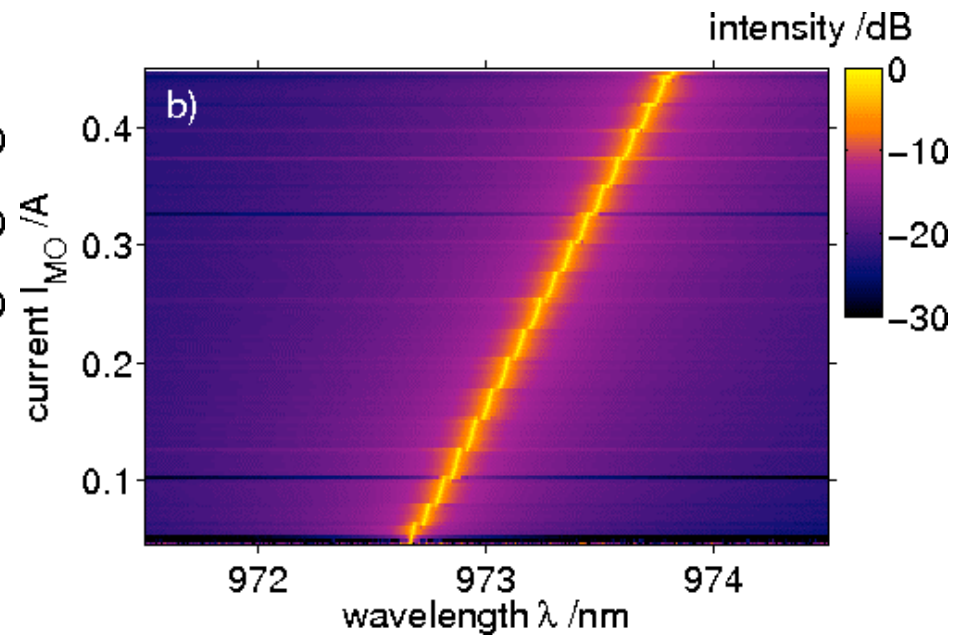
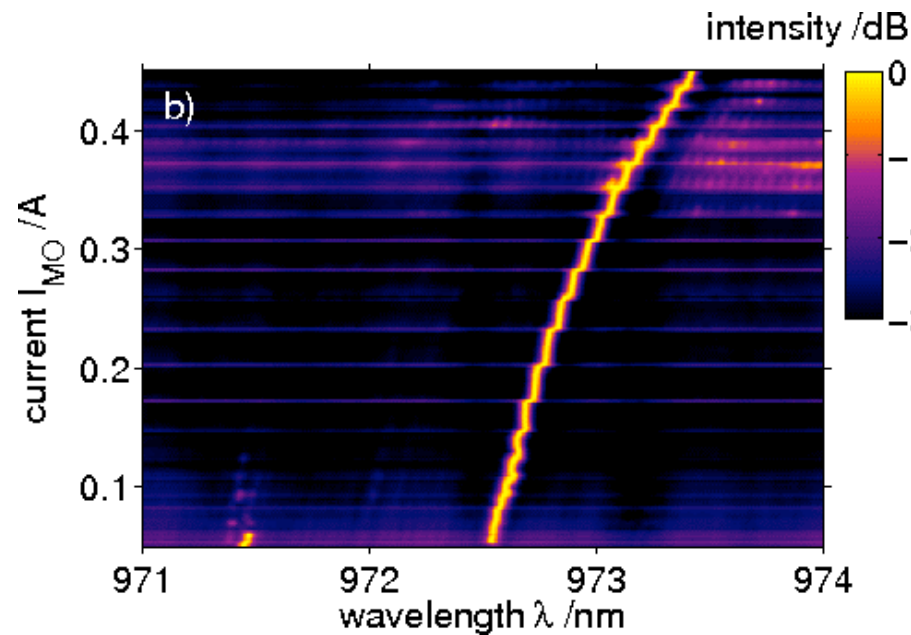


$$I_{MO} = 350 \text{ mA}$$

Optical Spectra vs I_{MO} :

Experiment:

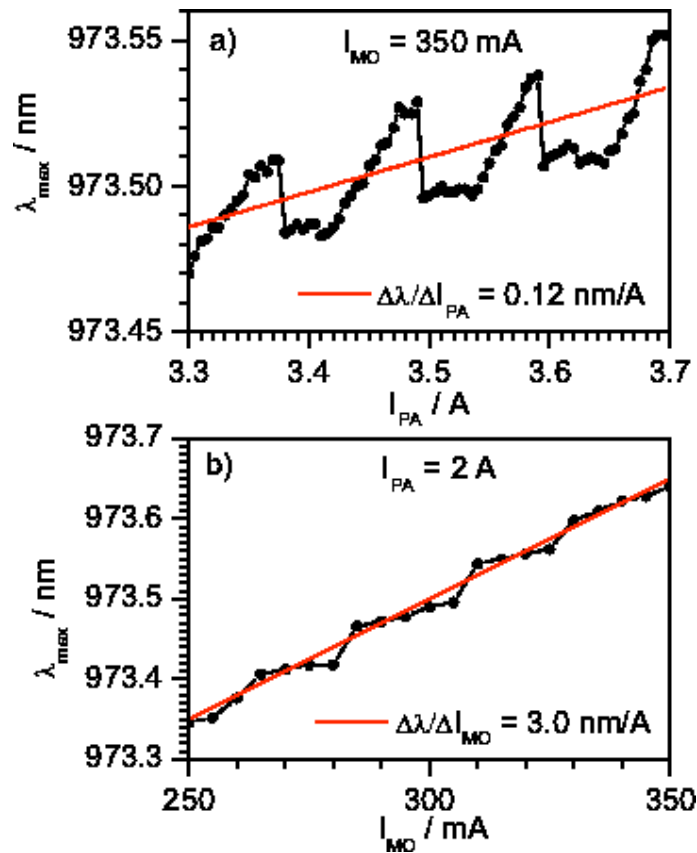
Simulation:



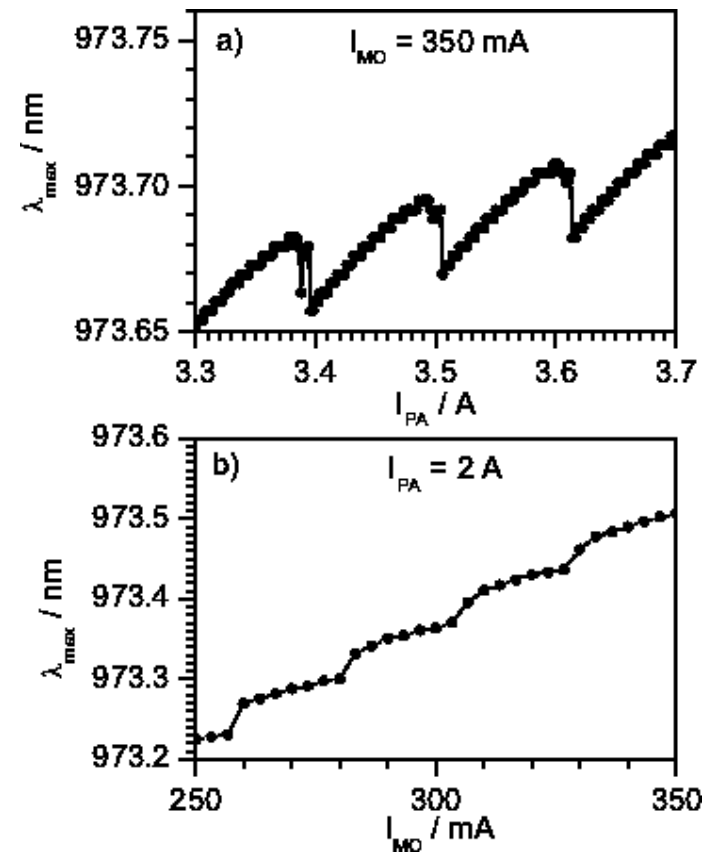
$$I_{PA} = 2 \text{ A}$$

Wavelength Shift

Experiment



Simulation



Summary

- Effective mathematical Model for MOPA's
- Heating effects parametrically included
- Detailed Comparison Experiment - Simulation
- Qualitative (longitudinal) understanding

Thank you for your Attention!